



<http://www.diva-portal.org>

This is the published version of a paper published in *Physical Review A. Atomic, Molecular, and Optical Physics*.

Citation for the original published paper (version of record):

Linusson, P., Hedin, L., Eland, J H., Squibb, R J., Mucke, M. et al. (2013)
Complete double valence photoionization study of the electron spectra of krypton.
Physical Review A. Atomic, Molecular, and Optical Physics, 88(2): 022510
<http://dx.doi.org/10.1103/PhysRevA.88.022510>

Access to the published version may require subscription.

N.B. When citing this work, cite the original published paper.

Permanent link to this version:

<http://urn.kb.se/resolve?urn=urn:nbn:se:uu:diva-207526>

Complete double valence photoionization study of the electron spectra of kryptonP. Linusson,¹ L. Hedin,² J. H. D. Eland,^{3,2} R. J. Squibb,² M. Mucke,² S. Zagorodskikh,² L. Karlsson,² and R. Feifel²¹*Department of Physics, Stockholm University, AlbaNova University Center, SE-106 91 Stockholm, Sweden*²*Department of Physics and Astronomy, Uppsala University, Box 516, SE-751 20 Uppsala, Sweden*³*Department of Chemistry, Physical and Theoretical Chemistry Laboratory, Oxford University, South Parks Road, Oxford OX1 3QZ, United Kingdom*

(Received 3 June 2013; published 21 August 2013)

Double photoionization spectra of Kr have been recorded using monochromatized synchrotron radiation of 88 eV photon energy and a versatile multielectron coincidence time-of-flight spectroscopy technique. The formation of the Kr^{2+} states of the lowest-energy configuration $4s^2 4p^4$ is partly direct, producing electron pairs with a continuous distribution, and partly indirect via superexcited singly ionized states. The superexcited Kr^+ states show strong and hitherto unexplained selectivity in branching to final Kr^{2+} states. Kr^{2+} states based on excited configurations are formed mainly by direct double photoionization.

DOI: [10.1103/PhysRevA.88.022510](https://doi.org/10.1103/PhysRevA.88.022510)

PACS number(s): 33.70.Ca, 33.80.Eh, 34.50.Gb

I. INTRODUCTION

Double ionization of atoms and molecules induced by single-photon absorption has been the subject of intense research for many years. A major focus of interest has been direct double ionization, meaning a simultaneous release of two electrons, in a process requiring strong electron correlation. The indirect double photoionization process, in which an intermediate singly ionized state is first formed (by a one-photon two-electron transition) and then decays by autoionization to a doubly charged ion, has been much less well studied. It is only with the advent of multicoincidence techniques that the second step, autoionization, in indirect double photoionization can be studied directly and the branching of particular monocationic levels to final dicationic states determined. Experimentally, the np^{-2} double ionization of the noble gases has been explored to a considerable extent (see, e.g., Refs. [1–4], and references therein). Double photoionization (DPI) leading to higher excited states of the dications, e.g., with a vacancy in the outermost s subshell, has received comparatively less attention. So far, neon has been thoroughly explored and the cross section for the formation of states related to the $2s^1 2p^5$ and $2s^0 2p^6$ configurations has been determined by both optical spectroscopy [5] and electron-electron coincidence spectroscopy [6]. Threshold photoelectron coincidence (TPEsCO) [7] studies have been carried out for photon energies covering the entire region below the core levels in neon and argon [8], as well as in xenon [9].

In this work, we present an electron-electron coincidence study of doubly ionized krypton atoms, using the time-of-flight magnetic bottle spectrometer technique with full energy multiplexing (Ref. [3], and references therein) at a photon energy of 88 eV. This is sufficient for the formation of dicationic states with energies up to and including the $4s^0 4p^6$ 1S state, but below the binding energy of the $3d$ core levels. Assignments of most of the lines, based upon data from high-resolution optical spectroscopy collected in the NIST database [10] and upon Auger electron spectroscopy studies [11,12], provide a much more complete characterization of the process than before. At the chosen photon energy, the Kr photoelectron spectrum [13] shows a strong (1D) nd 2S series of satellites, as well as a complicated satellite spectrum in

the binding-energy region above 40 eV. Their influence on the formation of the states derived from the $4p^4$ ground-state configuration is investigated in detail and shows remarkable selectivity in the autoionization branching.

II. EXPERIMENTAL DETAILS

The experiments were carried out using pulsed monochromatized synchrotron radiation as available at beam line U49/2 PGM-1 of the BESSY-II storage ring in Berlin, Germany. The storage ring was operated in single-bunch mode, providing light pulses at a rate of approximately 1.25 MHz. The magnetic bottle Time-Of-Flight (TOF) PhotoElectron-PhotoElectron COincidence (PEPECO) technique was used to measure the energies of all electrons liberated by single photons. Its principle has been reported earlier [4]. For the present studies, a 2.2-m-long magnetic bottle flight tube has been employed to determine the flight time of each of the two electrons ejected upon double ionization and detected in coincidence. Due to the long flight tube, near-zero kinetic-energy electrons have a typical flight time of $\sim 4 \mu\text{s}$, while those electrons with kinetic energies in excess of ~ 21 eV have TOFs shorter than the inter-light-pulse period (~ 800.5 ns). This means that the light pulses provided by the BESSY-II storage ring, even when operated in single bunch mode, do not provide an absolute time reference in particular for low-kinetic-energy electrons. For unambiguous determination of the TOF of the electrons, the light pulse frequency was therefore reduced to ~ 78 kHz by means of a mechanical chopper [14]. The electron counting rate was held at approximately 3 kHz and the photon energy resolution was below 10 meV. The sample gas was obtained commercially with a stated purity of better than 99%.

III. RESULTS AND DISCUSSION

Figure 1 shows an electron-electron coincidence map of krypton ionized by 88 eV photons, where the kinetic energy of the first arrival electron (fast electron) is plotted versus the kinetic energy of the second arrival electron (slow electron). The intensity in the map is distributed along the lines of the constant kinetic-energy sum of two electrons, which correspond to states of the dication, where a lower kinetic

structure of $4s4p^5\ ^3P$ is only partially resolved. At higher ionization energies, the spectrum becomes highly complex, as excited configurations of the type $4p^3nl$ give rise to many possible states. At around 62.7 eV ionization energy, an intense peak is observed. From studies of the M - NN Auger decay [15], it is known that strong configuration interaction (CI) between $4p^34d$ and $4s4p^5$ causes large redistribution of the intensity for the $J = 1$ levels, and therefore we attribute line 7 to be a $4p^34d\ ^1P_1$ satellite. Numerous other, less intense features are also observed in the energy region 52–65 eV. As individual peaks are not resolved, it is difficult to determine if there is a systematic behavior of the intensities. It appears (e.g., structure 3) that states with high J values ($J > 2$), whose intensity cannot be explained by CI with the $4s4p^5$ configuration, also acquire non-negligible intensities.

At ionization energies above 65 eV, five moderately strong peak structures are observed. Most of the intensity likely appears because of spectral artifacts originating from the double Auger (DA) decay of the $3d$ core holes [16,17] ionized by second- and third-order synchrotron light. While a triple ionization event in reality, the DA electrons appear in the two-electron coincidence data when the photoelectron is not detected. Their contribution to the intensity of the double-ionization spectrum can be estimated from triple electron coincidences, when the $3d$ photoelectron is detected, and is indicated in Fig. 2. However, part of feature 13 comes from double ionization and can be assigned to the $4s^04p^6\ ^1S$ state [10]. We note that this assignment disagrees with Ref. [15], but is supported by a recent calculation by Pernpointer *et al.* [18].

B. Double-ionization processes

The relative areas of the three outermost line structures reflect the probabilities of creating the 3P , 1D , and 1S states. The observed intensity ratios of 7.1:10.6:1, as obtained from the peak areas, are different compared to previous experimental observations using ionizing radiation from He II [4]. Part of the explanation is contributions to the intensity from indirect double ionization via intermediate singly ionized states. Obviously, some of them were not energetically accessible in the work of Ref. [4] and, in addition, those that were appear to be of less importance at the present photon energy.

In Fig. 3(a), the low-kinetic-energy part of the electron-energy distribution associated with the formation of the $4p^4\ ^3P_{2,1,0}$ states is shown. Line positions and relative intensities of the features observed are given in Table II. The observed peaks arise because of autoionization to the 3P states; for comparison with photoelectron spectroscopy (PES) studies, one of the two energy scales represents the electron kinetic energy plus the formation energy for the 3P_2 ground state of the dication. Here and subsequently, we assume an energy of 38.3594 eV for the 3P_2 state according to the NIST database [10]. At the present photon energy of 88 eV, the most intense satellite features in the Kr photoelectron spectrum in the binding energy region 38–40 eV are attributed to the $(^1D)nd\ ^2S$ states [13,19,20]. In the autoionization spectrum of Fig. 3(a), this is manifested by a series of strong peaks (3,7,11,13,15–18) corresponding to $(^1D)nd\ ^2S \rightarrow ^3P_2$ transitions. A copy of this series (peaks 1,5,6,8–10) can also be discerned for $n = 8–13$, which differs for each n in energy to the original very nearly by the $^3P_{2,1}$ level

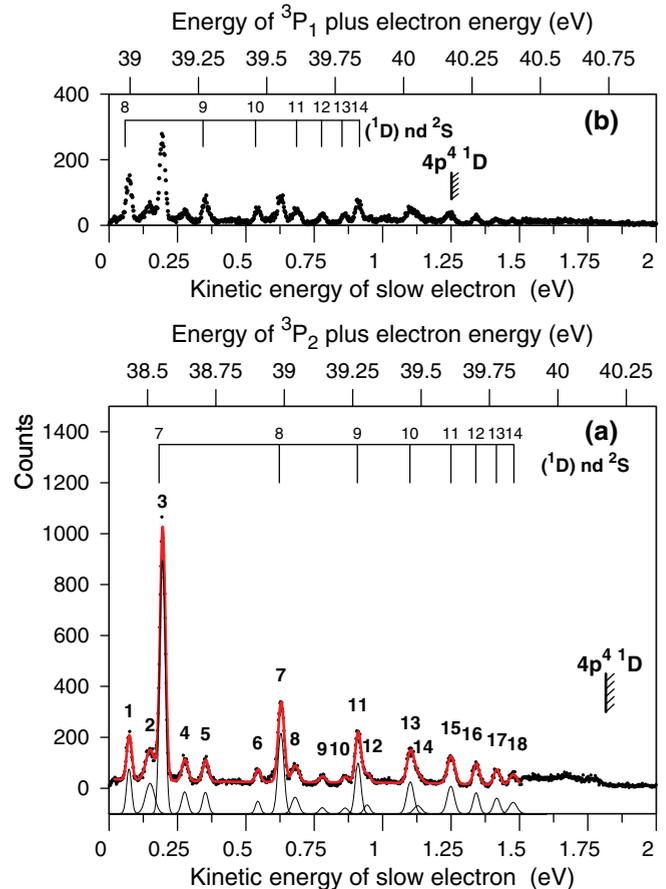


FIG. 3. (Color online) Spectra of electrons released in the formation of the $4p^4\ ^3P$ states obtained from different selections on the kinetic-energy sum of the two electrons. (a) The entire peak (cf. Fig. 2) corresponding to the 3P states has been selected. Dots show experimental results and solid lines represent a fit to Gaussian peak shapes. The top scale is the electron kinetic energy plus the energy of the 3P_2 ground state of the dication relative to neutral Kr. (b) A subset of the data in an energy region (indicated in Fig. 2) corresponding to the $^3P_{1,0}$ final states has been selected. The step size is 2 meV in both (a) and (b).

splitting in the dication. The relative intensity of the repeated series is enhanced in Fig. 3(b), where we have selected a subset of the coincidence data according to an energy which corresponds to the formation of the $^3P_{1,0}$ final states. We note that overlap with the decay to the 3P_2 state occurs, in particular, for $n = 11, 14$. In addition, a weak feature at a binding energy approximately corresponding to peak 1 in Fig. 3(a) was observed in a photoelectron spectrum recorded at 68.5 eV by Caló *et al.* [20]. However, it was not reported in the 125 eV study by Alitalo *et al.* [19], and we note that peak 5 in Fig. 3(a) has no correspondence in the photoelectron spectra of either study. In view of these observations, we assign the repeated series as dominated by $(^1D)nd\ ^2S \rightarrow ^3P_1$ autoionization and estimate a branching ratio of the decay of the $(^1D)nd\ ^2S$ ($n = 8–10$) singly ionized states to the $^3P\ J = 2; 1; 0$ levels of the dication to be approximately 2.2; 1; ~ 0 . Further experimental effort is required to deduce more accurate

branching ratios, as it is not clear from the present results whether this ratio changes for higher n .

It is interesting to compare the present results with the $np^4(^1D)nd^2S$ autoionizing resonances in the halogens, particularly Br which is isoelectronic to Kr^+ . In strict LS coupling, the autoionization of the $(^1D)nd^2S$ Rydberg series in the halogens to the 3P states is “forbidden,” but occurs as a consequence of spin-orbit interaction [21] as the $(^1D)nd^2S$ resonance states can contain small admixtures of $(^1D)nd^2P_{1/2}$ [22] and the “target” 3P_2 and 1D_2 cores are coupled [23]. Autoionization to the 3P_2 final state was found experimentally to be dominant for the $(^1D)nd^2S$ resonances in Cl and Br [24,25], and only weak transitions to the $J = 1, 0$ levels were observed. The apparently much smaller $^3P_2/3P_1$ branching ratio observed here for the decay of the Kr^+ states compared to Br may, to some extent, be explained by the increased effective charge of the nucleus, increasing spin-orbit effects. Another obvious difference is the presence of an escaping photoelectron in the present case so that coupling between the continuum channels should be considered. However, estimating the importance of such couplings will require theoretical investigations which

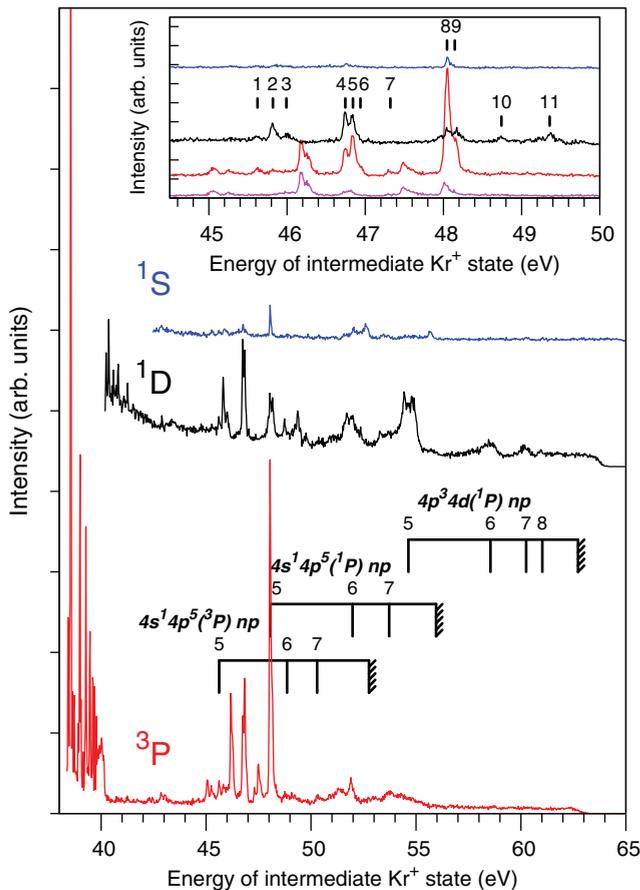


FIG. 4. (Color online) Spectrum of electrons released in the formation of the $4p^4$ 3P , 1D , and 1S states of the Kr dication. For the lowest trace in the inset, which is essentially a zoomed-in version of the main figure, a final-state energy region corresponding to the formation of the $^3P_{1,0}$ states (cf. Fig. 3) has been selected. The energy scale is the ionization energy of the indicated state plus the electron kinetic energy. The intensity scale is common, but traces have been offset vertically.

TABLE II. Line positions and intensities observed in the low kinetic energy (<2 eV) formation spectrum of the $4p^4$ 3P states. The last column denotes binding energies (BE). The error in relative energies is estimated to be within 10 meV. The relative intensities (Rel. Int.) are expected to be accurate within 10% unless otherwise stated (in form of giving the actual deviations as numbers in parentheses).

Line	E_{kin} (eV)	Rel. Int.	Decay	BE (eV)
1	0.073	16.3	$^1D\ 8d\ ^2S \rightarrow ^3P_1$	38.996
2	0.150	20.0	$^1S\ 6p \rightarrow ^3P_2$	38.509
3	0.195	100	$^1D\ 7d\ ^2S \rightarrow ^3P_2$	38.555
4	0.276	9.7	$^1D\ 8d\ ^2D_{3/2} \rightarrow ^3P_2$	38.636
5	0.352	8.9	$^1D\ 9d\ ^2S \rightarrow ^3P_1$	39.275
6	0.544	4.7	$^1D\ 10d\ ^2S \rightarrow ^3P_1$	39.468
7	0.628	35.9	$^1D\ 8d\ ^2S \rightarrow ^3P_2$	38.988
8	0.680	8.6	$^1D\ 11d\ ^2S \rightarrow ^3P_1$	39.604
9	0.780	2.4(5)	$^1D\ 12d\ ^2S \rightarrow ^3P_1$	39.703
10	0.863	2.3(5)	$^1D\ 13d\ ^2S \rightarrow ^3P_1$	39.787
11	0.910	21.8	$^1D\ 9d\ ^2S \rightarrow ^3P_2$	39.269
12	0.943	3.9	$^1D\ 8d\ ^2D_{3/2} \rightarrow ^3P_2$	39.291
13	1.101	17	$^1D\ 10d\ ^2S \rightarrow ^3P_2$	39.460
14	1.13	5(1)	$^1D\ 10d\ ^2D_{3/2} \rightarrow ^3P_{2,1}$	39.49
15	1.249	15	$^1D\ 11d\ ^2S \rightarrow ^3P_2$	39.609
16	1.342	9	$^1D\ 12d\ ^2S \rightarrow ^3P_2$	39.701
17	1.417	8	$^1D\ 13d\ ^2S \rightarrow ^3P_2$	39.776
18	1.477	7	$^1D\ 14d\ ^2S \rightarrow ^3P_2$	39.836

are beyond the scope of the present work. We note, as already pointed out in Ref. [4], that a weak decay to the 3P_0 final state can be more readily explained, as the 3P_0 ϵd continua are not allowed for a $J = 1/2$ state and hence a change of angular momentum for the Rydberg electron would be required.

A number of weaker features, peaks 2 and 4, and the partly resolved peaks 12 and 14 are also seen in Fig. 3(a). Assignments of them are not clear cut, as many states may contribute [19,26], but in the study of Caló *et al.* [20] the previous assignment [19] of a $^1S\ 6p$ shakeup state at a binding energy close to that of peak 2 was confirmed, and new assignments of $(^1D)nd^2D_{3/2}$ states close to that of peaks 12 and 14 were made. More information on states near the apparent binding energy of peak 4 was, unfortunately, not reported in the work of Caló *et al.* [20], but we note that it has an apparent binding energy close to the $(^1D)8d^2D_{3/2}$ state, which appears strongly in threshold photoelectron spectra [27,28] and (probably) at the He II photon energies [4]. Hence, we tentatively assign it to that state.

In Fig. 4, the distributions for the formation of the states related to the $4p^4$ configuration are shown on a single-ionization energy scale. The photoelectron spectrum of Kr in the ionization energy region 42–65 eV was included already in the study by Kikas *et al.* [13] and is very complex. The many overlapping spectral features make assignments difficult. However, below 50 eV ionization energy, some clear lines can be distinguished, which are summarized in Table III along with estimates of the branching ratios of the decay. In the distribution for the 1D state, some sharp peaks are also observed for the 40–42 eV binding-energy range, but detailed

TABLE III. Line positions, relative intensities, and autoionization branching ratios in the binding-energy region 44–50 eV. Line numbers refer to Fig. 4.

Line	BE (eV)	Assignment (Ref. [13])	Rel. Int. (arb. units)	Branching (%)			
				3P_2	$^3P_{1,0}$	1D_2	1S_0
1	45.6	$4s4p^5(^3P)5p$	7	44	56		
2	45.8		18		28	72	
3	46.0		12		18	82	
4	46.7	$4s4p^5(^1P)5s$	64	29	37	31	4
5	46.8	$4s4p^5(^3P)4d$	61	48	22	30	
6	46.9		10	49	51		
7	47.3		2	100			
8	48.0	$4s4p^5(^1P)5p$	100	76	10	8	6
9	48.1		50	76	7	17	
10	48.7		3			100	
11	49.4		12			100	

analysis suggests that they should be attributed rather to the 3P state, and appear here due to the imperfect final-state selection.

The $4s4p^5np$ states have been studied extensively in the spectator Auger decay of the $3d^{-1}np$ optical resonances [15,29–31], and approximate positions are indicated in Fig. 4. Such states can arise due to monopole shakeup accompanying the $4s$ ionization, which is expected to dominate at high energies. However, in analyzing the corresponding satellite spectrum of argon at 100 eV photon energy, Combet-Farnoux *et al.* [32] found that electron correlations, in particular in the ground state, were more important in explaining the intensity of the satellites. A similar case can be expected for the present system and a detailed understanding of the distributions shown in Fig. 4 requires theoretical input. The single-configuration notation used here should be regarded as convenient labels; few of the states are expected to be dominated by single-configuration weights [31].

Regardless of the formation process, the selectivity in autoionization of the satellites to the states of the $4p^4$ configuration is striking. The $4s4p^54d$, $4s4p^55s$, and, in particular, $4s4p^55p$ satellites all favor autoionization to the $4p^4\ ^3P_{2,1}$ states. We note that this differs from the Ar case, where autoionization to the 1D final state dominates for the most intense satellites [32]. At higher ionization energies, most of the intensity is observed as autoionization to the 1D state. The most prominent example is the series of broad and asymmetric peaks leading up to the double-ionization limit of the intense $4p^34d\ ^1P$ satellite (cf. Fig. 2), which is clearly observed in the $4p^4\ ^1D$ distribution but is very weak in the 3P and 1S distributions. Peaks at similar energies do, however, also appear in the electron distributions related to the $4s4p^5$ final states (not shown).

The intensity of the direct DPI channel can be estimated from the data by counting the continuum parts of the energy distributions, as has been done in previous works [4,6]. We follow the work of Ref. [6] and estimate the continuum parts by fitting a hyperbolic function to a number of points in the distributions. In the present case, this approach appears reasonable for the $4p^4\ ^3P$ and 1S states, as regions comparatively

free of resonance structure are apparent at several points in this energy range, but considerable difficulties are met for the 1D state, as essentially the whole available energy range shows some involvement of indirect channels and interference can be present. Keeping the limitations of the approach in mind (see Ref. [6]), we estimate the continuum parts to account for 40%, 50%, and 70% of the intensity of the $4p^4\ ^3P$, 1D , and 1S states, respectively. In a similar way, the continuum parts of the distributions related to the $4s4p^5\ ^3P$ and 1P states account for approximately 60% and 80%, respectively, of the intensity. The distributions of higher-lying dicationic states appear essentially free of peak structures.

IV. CONCLUSIONS

The complete double valence photoionization spectrum of krypton has been investigated using the magnetic bottle TOF-PEPECO method. In the formation of the 3P , 1D , and 1S states of the $4p^4$ configuration, indirect double ionization via intermediate singly ionized states account for at least half of the total intensity. The autoionization processes show strong selectivity, which cannot be explained on the basis of known propensity rules and thus presents a challenge to existing theory.

ACKNOWLEDGMENTS

This work has been financially supported by the Swedish Research Council (VR), the Göran Gustafsson Foundation (UU/KTH), the Knut and Alice Wallenberg Foundation, and the Carl Tryggers Foundation, Sweden. We would like to warmly acknowledge the support by the staff and colleagues at BESSY-II, Berlin. This work was also supported by the European Community - Research Infrastructure Action under the FP6 “Structuring the European Research Area” Programme (through the Integrated Infrastructure Initiative “Integrating Activity on Synchrotron and Free Electron Laser Science” - Contract No. R II 3-CT-2004-506008).

- [1] P. Bolognesi, G. C. King, and L. Avaldi, *Rad. Phys. Chem.* **70**, 207 (2004).
- [2] G. C. King and L. Avaldi, *J. Phys. B: At. Mol. Opt. Phys.* **33**, R215 (2000).
- [3] J. H. D. Eland, *Adv. Chem. Phys.* **141**, 103 (2009).
- [4] J. H. D. Eland, O. Vieuxmaire, T. Kinugawa, P. Lablanquie, R. I. Hall, and F. Penent, *Phys. Rev. Lett.* **90**, 053003 (2003).
- [5] K. H. Schartner, G. Mentzel, B. Magel, B. Mobus, A. Ehresmann, F. Vollweiler, and H. Schmoranzler, *J. Phys. B: At. Mol. Opt. Phys.* **26**, L445 (1993).
- [6] T. Kaneyasu, Y. Hikosaka, E. Shigemasa, F. Penent, P. Lablanquie, T. Aoto, and K. Ito, *Phys. Rev. A* **76**, 012717 (2007).
- [7] R. I. Hall, A. McConkey, K. Ellis, G. Dawber, L. Avaldi, M. A. MacDonald, and G. C. King, *Meas. Sci. Technol.* **3**, 316 (1992).
- [8] L. Avaldi, G. Dawber, N. Gulley, H. Rojas, G. C. King, R. Hall, M. Stuhec, and M. Zitnik, *J. Phys. B: At. Mol. Opt. Phys.* **30**, 5197 (1997).
- [9] P. Bolognesi, S. J. Cavanagh, L. Avaldi, R. Camilloni, M. Zitnik, M. Stuhec, and G. C. King, *J. Phys. B* **33**, 4723 (2000).
- [10] Y. Ralchenko, A. Kramida, J. Reader, and NIST ASD Team (2010), *NIST Atomic Spectra Database (version 4.0.1)*, <http://physics.nist.gov/asd> (National Institute of Standards and Technology, Gaithersburg, MD, 2013).
- [11] H. Pulkkinen, S. Aksela, O.-P. Sairanen, A. Hiltunen, and H. Aksela, *J. Phys. B: At. Mol. Opt. Phys.* **29**, 3033 (1996).
- [12] H. Aksela, S. Aksela, and H. Pulkkinen, *Phys. Rev. A* **30**, 2456 (1984).
- [13] A. Kikas, S. Osborne, A. Ausmees, S. Svensson, O.-P. Sairanen, and S. Aksela, *J. Elec. Spec. Rel. Phen.* **77**, 241 (1996).
- [14] S. Plogmaker, P. Linusson, J. H. D. Eland, N. Baker, E. M. J. Johansson, H. Rensmo, R. Feifel, and H. Siegbahn, *Rev. Sci. Instrum.* **83**, 013115 (2012).
- [15] H. Aksela, S. Aksela, H. Pulkkinen, G. M. Bancroft, and K. H. Tan, *Phys. Rev. A* **33**, 3876 (1986).
- [16] E. Andersson, S. Fritzsche, P. Linusson, L. Hedin, J. H. D. Eland, J.-E. Rubensson, L. Karlsson, and R. Feifel, *Phys. Rev. A* **82**, 043418 (2010).
- [17] J. Palaudoux, P. Lablanquie, L. Andric, K. Ito, E. Shigemasa, J.H.D. Eland, V. Jonauskas, S. Kucas, R. Karazija, and F. Penent, *Phys. Rev. A* **82**, 043419 (2010).
- [18] M. Pernpointner, J. P. Zobel, and N. V. Kryzhevoi, *Phys. Rev. A* **85**, 012505 (2012).
- [19] S. Alitalo, A. Kivimäki, T. Matila, K. Vaarala, H. Aksela, and S. Aksela, *J. Elec. Spec. Rel. Phen.* **114–116**, 141 (2001).
- [20] A. Caló, S. Atanassova, R. Sankari, A. Kivimäki, H. Aksela, and S. Aksela, *J. Phys. B: At. Mol. Opt. Phys.* **39**, 4169 (2006).
- [21] J. Berkowitz, in *Advances in Chemical Physics*, edited by I. Prigogine and Stuart A. Rice (Wiley, New York, 2007), Vol. 72, pp. 1–36.
- [22] J. E. Hansen, R. D. Cowan, S. L. Carter, and H. P. Kelly, *Phys. Rev. A* **30**, 1540 (1984).
- [23] F. Robicheaux and C. H. Greene, *Phys. Rev. A* **46**, 3821 (1992).
- [24] S. Benzaid, A. Menzel, J. Jiménez-Mier, S. J. Schaphorst, M. O. Krause, and C. D. Caldwell, *Phys. Rev. A* **54**, R2537 (1996).
- [25] S. Benzaid, M. O. Krause, A. Menzel, and C. D. Caldwell, *Phys. Rev. A* **57**, 4420 (1998).
- [26] J. Jauhiainen, H. Aksela, O.-P. Sairanen, E. Nömmiste, and S. Aksela, *J. Phys. B: At. Mol. Opt. Phys.* **29**, 3385 (1996).
- [27] M. Hochlaf, H. Kjeldsen, F. Penent, R. I. Hall, P. Lablanquie, M. Lavollé, and J. H. D. Eland, *Can. J. Phys.* **74**, 856 (1996).
- [28] H. Yoshii, T. Aoto, Y. Morioka, and T. Hayaishi, *J. Phys. B: At. Mol. Opt. Phys.* **40**, 2765 (2007).
- [29] H. Aksela, S. Aksela, A. Mäntykenttä, J. Tulkki, E. Shigemasa, A. Yagishita, and Y. Furusawa, *Phys. Scr.* **T41**, 113 (1992).
- [30] H. Aksela, J. Jauhiainen, E. Kukkk, E. Nömmiste, S. Aksela, and J. Tulkki, *Phys. Rev. A* **53**, 290 (1996).
- [31] J. Mursu, J. Jauhiainen, H. Aksela, and S. Aksela, *J. Phys. B: At. Mol. Opt. Phys.* **31**, 1973 (1998).
- [32] F. Combet-Farnoux, P. Lablanquie, J. Mazeau, and A. Huetz, *J. Phys. B: At. Mol. Opt. Phys.* **33**, 1597 (2000).