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# LOW FREQUENCY DIELECTRIC PROPERTIES OF THREE BENTONITES

## AT DIFFERENT ADSORBED WATER STATES

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9  
10 KEYWORDS: clay mineral, layer silicate, smectite, bound water, permittivity, equivalent  
11 circuit

### 12 13 ABSTRACT

14 Three bentonites of varying smectite content were investigated by dielectric spectroscopy in  
15 the frequency range  $10^{-4}$  to  $10^6$  Hz after storage at well-defined humidities. The identification  
16 of relaxation processes from complex permittivity measurements was difficult, since  
17 conductivity effects were superimposed on the underlying relaxations. Relaxation peaks  
18 revealed by the dissipation factor indicated the occurrence of interfacial processes between  
19  $10^2$  and  $10^6$  Hz. The intensity of the polarization of the electrochemical double-layer at the  
20 clay-water interface was promoted by increasing water content and was shifted to higher  
21 frequencies the higher the water content in the bentonites. Below  $\sim 1$  Hz electrode  
22 polarization (EP) was shown to be a participating process with capacitance values ranging  
23 from  $0.6 \cdot 10^{-3}$  to  $7.3 \cdot 10^{-3}$  F due to the accumulated charges. An equivalent circuit model was  
24 introduced that successfully described the low frequency dielectric behaviour of bentonites at  
25 low moisture levels. An included series R-CPE connection was used to describe the double-  
26 layer relaxation. At water contents up to 17% the bulk resistivity was mainly influenced by  
27 smectite content and cation exchange capacity, whereas at water contents of  $\geq 19\%$   
28 interlayer occupation and hydration state became more important.

## 30 INTRODUCTION

31 Bentonites are natural materials occurring on surface and subsurface deposits around the  
32 world <sup>[1]</sup>. They consist of major amounts of the swellable clay mineral montmorillonite or other  
33 minerals of the smectite group. Since even the swellable 2:1 layer silicate montmorillonite  
34 widely varies in composition and structure <sup>[2]</sup>, a detailed knowledge about their mineralogical  
35 features is substantial to understand the influence on the macroscopic behaviour of the  
36 bentonites <sup>[3-5]</sup>.

37 Due to specific properties of smectites, for instance high sorption capacity and their ability to  
38 reversibly exchange cations, smectite-rich materials, such as bentonites, are widely used for  
39 industrial and environmental applications. They have gained an increasing importance in the  
40 field of waste disposal of hazardous materials, like radioactive substances, but are also –  
41 along with other clays – used as fillers and major components in cosmetics and  
42 pharmaceuticals <sup>[6, 7]</sup>. In agriculture, high clay content together with the presence of 2:1 clay  
43 minerals is associated with high soil fertility. Processes due to clay mineral intrinsic  
44 properties, such as water retention and provision of exchangeable sites for the essential  
45 nutrients in plant growth, are predominant features increasing soil fertility <sup>[1]</sup>.

46 For many environmental and industrial applications it is crucial to accurately adjust the water  
47 content of the mineral mixtures. Hence, the precise determination of water content is  
48 necessary for instance to ensure high quality products and to assess the stability of the  
49 geotechnical barriers on waste disposal sites. However, the water binding mechanisms are  
50 complex in clay minerals <sup>[8-11]</sup>. Swellable clay minerals bind water on their outer and on their  
51 inner surface upon hydration, whereas the interlayer cations undergo stepwise hydration,  
52 which is influenced by valence and radii of the cations <sup>[12]</sup>. With increasing relative humidity  
53 an increasing amount of water molecules is associated to the interlayer cations that can be  
54 described by different hydration states ranging from 0W to 3W <sup>[13]</sup>.

55 Several techniques were established for the determination of water content <sup>[14]</sup>, among which  
56 static heating at 105 °C until weight equilibrium is commonly used and regularly applied as a  
57 reference method <sup>[15]</sup>. Beside heating, indirect physical methods like tensiometric techniques,

58 geophysical methods (e.g. electrical conductivity technique), nuclear approaches (e.g.  
59 neutron-scattering technique) as well as several high frequency electrical techniques <sup>[16]</sup> and  
60 remote sensing techniques <sup>[17]</sup> are in use <sup>[18, 19]</sup>. Each technique is characterized by individual  
61 features, such as accuracy, spatial resolution, time consumption and cost, and further holds  
62 different requirements on sample amount and preparation. The preciseness of the different  
63 water content measurement approaches may vary distinctively, especially when determining  
64 the water content of materials containing substantial amounts of swellable clay minerals. In  
65 many applications, such as the assessment of geotechnical barriers in landfill sites, in-situ  
66 and non-destructive water content determination is required.

67 Thus, moisture measurement by means of dielectric spectroscopy, which uses the  
68 relationship between the real relative permittivity and the volumetric water content of a  
69 sample, is a promising technique, and it has been shown that both the state of water (bound  
70 or free) <sup>[20]</sup> as well as the number of dissociated water molecules carrying charges in a  
71 system <sup>[21-23]</sup> can be studied with this technique. Moreover, it is not restricted by safety  
72 regulations in contrast to nuclear approaches <sup>[19, 24]</sup>. However, recent work shows that the  
73 regularly used frequency of about 1 GHz seems to be unsuitable for water content  
74 determination with dielectric spectroscopy when the studied materials contain a major  
75 amount of swellable clay minerals <sup>[25]</sup>.

76 A wide range of studies was carried out regarding the determination of volumetric water  
77 content of soils by dielectric spectroscopy <sup>[26-30]</sup> and they showed that a number of  
78 parameters, such as bulk density <sup>[31]</sup> and iron content <sup>[32]</sup> influences the results in laboratory  
79 and in the field. Moreover, swelling clays <sup>[33, 34]</sup> as well as non-swelling clays <sup>[35]</sup> were  
80 investigated by dielectric spectroscopy over a wide moisture range to study the clay specific  
81 behaviour such as flocculation and gel formation and the role of cations as mobile charge  
82 carriers.

83 Most of the dielectric studies performed on clays were conducted at high moisture levels, on  
84 suspensions or pastes <sup>[36-41]</sup>. However, in many of the industrial and environmental

85 applications where the moisture is required to be monitored, such as for sealing layers in  
86 hazardous waste disposal sites, clays are present at much lower water contents.

87 Our own investigations on swellable and non-swellable 2:1 layer silicates showed a higher  
88 sensitivity of low-frequency measurements ( $10^{-4}$  ...  $10^6$  Hz) compared to high-frequency  
89 measurements ( $2.0 \cdot 10^8$  ...  $1.1 \cdot 10^9$  Hz) to display differences in the adsorbed water state [25].  
90 Still, low frequency dielectric analysis of clay-water-systems has commonly been conducted  
91 on clay suspensions or clay-water-electrolyte systems [34, 40-42]. In contrast to mid and high  
92 frequency measurements, within the low-frequency range further polarization mechanisms,  
93 e.g. interfacial polarization, double-layer polarization and electrode polarization, may  
94 influence the dielectric spectra and superimpose on the bound water effect. A four terminal  
95 probe was suggested by Carrier and Soga [43] as a preliminary method to reduce electrode  
96 polarization effects within the low frequency range. Furthermore, some experimental and  
97 mathematical approaches were established to correct dielectric data for the contribution of  
98 electrode polarization [44-47].

99 Previous equivalent circuit fitting of clay-water-systems similarly concentrated on clay  
100 suspensions [40, 42, 48]. Dielectric studies on clays at low water content states and their  
101 assessment by equivalent circuit fitting is rarely described in the literature.

102 The aim of the current study is to investigate the dielectric behaviour of three bentonites with  
103 varying smectite content within the low-frequency range. The materials are studied at  
104 different adsorbed water states rather than as clay suspensions or clay pastes. Additionally,  
105 equivalent circuit fitting was implemented in order to evaluate the influence of the attributing  
106 polarization mechanisms as a function of varying humidity.

107  
108

## 109 MATERIALS

110 Dielectric properties of three bentonites were studied at different hydration states in the low-  
111 frequency range. The bentonite "Volclay" (ben\_Vol), "Calcigel" (ben\_Cal) and bentonite "P"  
112 (ben\_P) were provided by Süd-Chemie AG, Germany. Basic physical and mineralogical  
113 properties, such as mineralogical composition, cation exchange capacity (CEC), type of

114 exchangeable cations, grain size distribution as well as surface properties such as outer  
115 specific surface area ( $A_{s, out}$ ), total specific surface area ( $A_{s, total}$ ), and pore size distribution were  
116 determined. Hereby, the mineralogical composition was determined by means of X-ray  
117 diffraction analysis and the grain size distribution by sieving and sedimentation analysis. The  
118 CEC was determined photometrically according to the Cu(II)-Triethylenetetramine method <sup>[49]</sup>  
119 and the type of exchangeable cation of the supernatant from the copper-complex exchange  
120 reaction by inductively coupled plasma - optical emission spectroscopy (ICP-OES). Surface  
121 properties were investigated by nitrogen and water vapour adsorption isotherms, whereas  
122 the specific surface area were determined according to Brunauer-Emmett-Teller theory <sup>[50]</sup>  
123 and the pore size distribution according to the non-local density functional theory <sup>[51]</sup> using  
124 the adsorption branch model for cylindrical silica pores. A detailed description of the related  
125 experiments can be found in <sup>[52]</sup>.

126 The fine grained powders were stored under defined relative humidity (r.h.) conditions and  
127 prevented from previous heating or drying before the hydration experiments.

128 Among the studied bentonites, the ben\_P exhibits the lowest amount of impurities, consisting  
129 of about 98% of swellable minerals belonging to the smectite group (Table 1). Ben\_Vol is  
130 characterized by the second highest smectite content being about 78%. The material  
131 furthermore contains about 10% feldspars, 8% mica/illite, 4% quartz and traces of gypsum.  
132 The third material, ben\_Cal, exhibits the smallest smectite content (about 65%) and the  
133 highest amount of impurities. Major phases beside the smectites are non-swellable 2:1 layer  
134 silicates, such as mica/illite (17%) and 7% of quartz. The clay content of the materials, i.e.  
135 the fraction of  $< 2\mu\text{m}$ , decreases in the order ben\_P > ben\_Vol > ben\_Cal (Table 1).

136 The different mineralogical composition and particle size distribution of the bentonites is  
137 further displayed by varying exchange and surface properties. The CEC for instance is  
138 connected to the amount of swellable phases and their layer charge. The ben\_P exhibits the  
139 highest CEC (119 meq/100g) due to its high smectite content and the high layer charge of  
140 the smectite. The ben\_Vol exhibits an intermediate CEC (83 meq/100g) and the ben\_Cal the  
141 lowest CEC of the three bentonites (63 meq/100g). Since the bentonites are natural bulk

142 materials their interlayer occupation is heterogeneous (Table 1). The interlayer of the  
143 smectite in ben\_P is equally occupied by monovalent (51%) and divalent cations (49%),  
144 whereas the ben\_Vol is slightly dominated by monovalent cations (64%) and the ben\_Cal is  
145 clearly dominated by divalent cations (89%).

146 The external surface area ( $A_{s, out}$ ) is primarily influenced by grain size and porosity of a  
147 material. Since ben\_P displays the highest clay content (79%) of the three bentonites, a high  
148  $A_{s, out}$  of 105 m<sup>2</sup>/g is in agreement with the grain size distribution. However, ben\_Vol with a  
149 similarly high clay content exhibits a distinctively smaller  $A_{s, out}$  of about 33 m<sup>2</sup>/g, whereas  
150 ben\_Cal, which is the coarsest of the materials with a silt content of about 50%, exhibits a  
151 distinctively higher  $A_{s, out}$  of about 72 m<sup>2</sup>/g. Accordingly, nitrogen adsorption experiments and  
152 analysis by means of non-local density functional theory revealed that ben\_Cal is  
153 characterized by higher cumulated surface area of the pores than ben\_Vol (Table 1).

154 The total specific surface area ( $A_{s, total}$ ), which characterizes the sum of the external surface  
155 area as well as the interlayer surface, is strongly influenced by the hydration properties of the  
156 material, such as accessibility of the interlayer displayed by CEC as well as the type of  
157 interlayer cation. Consequently, the  $A_{s, total}$  decreases with decreasing CEC in the order  
158 ben\_P > ben\_Vol > ben\_Cal.



159

160 Table 1: Basic physical and mineralogical properties including mineralogical composition, type of interlayer cation,  
 161 outer specific surface area ( $A_{s, out}$ ) and total specific surface area ( $A_{s, total}$ ) as well as cumulative surface area of  
 162 micropores ( $A_{s, micro}$ ) and mesopores ( $A_{s, meso}$ ).

	Unit	Ben_Cal	Ben_Vol	Ben_P
Particle density	[g/cm <sup>3</sup> ]	2.66	2.69	2.65
CEC	[meq/100g]	63	83	119
Exchangeable cations	Na	9	62	50
	K	2	2	1
	Ca	62	30	40
	Mg	27	6	9
Grain size	Sand	3	3	0
	Silt	50	20	21
	Clay	47	77	79
$A_{s, out}$		72	33	105
$A_{s, total}$		289	378	417
$A_{s, micro}$	d < 2 nm	49	19	75
$A_{s, meso}$	d < 50 nm	78	35	117
Mineral	Calcite	3	-	-
	Dolomite	5	-	-
	Gypsum	-	< 1	-
	Kaolinite	3	-	-
	Mica/illite	17	8	1
	K-Feldspars	-	6	-
	Plagioclases	-	4	-
	Quartz	7	4	1
	Diocahedral smectite	65	78	98

163

164

## 165 METHODS

### 166 *Pre-treatment*

167 The powdered materials as received were stored in desiccators above saturated salt  
 168 solutions until equilibration to specific r.h.. The salt solutions used for humidity equilibration  
 169 were LiCl for 11% r.h., K<sub>2</sub>CO<sub>3</sub> for 43% r.h. and KNO<sub>3</sub> for 93% r.h..

170

### 171 *Determination of hydration state*

172 The main hydration state of the smectites was evaluated by determining the predominant  
 173 number of hydration shells around the interlayer cations. Since the position of the d001 basal  
 174 reflections of the smectites is a function of the extent of interlayer cation hydration, the

175 stepwise expansion of the smectites upon hydration, which is caused by intercalation of 0, 1,  
176 2 or 3 planes of water within the interlayer, can be determined from the shift of the  $d001$  with  
177 varying water content of the sample <sup>[53]</sup>. Hereby, powder X-ray diffraction (XRD) specimen of  
178 the equilibrated samples were prepared and measured using a diffractometer with  
179 monochromator, apertures of 1 mm and 0.1 mm, and  $\text{CuK}\alpha$  radiation in a range of  
180 5 to 35°  $2\theta$  with a step size of 0.03°  $2\theta/3$  s (D5000 diffractometer, Siemens, Germany) and  
181 the position of the  $d001$  was determined. Different hydration states are characterized for  
182 smectites as shown in Table 2.

183

184 Table 2: Hydration states of smectites, according to Ferrage *et al.* <sup>[53]</sup>.

Hydration state	Abbreviation	$d001$ [Å]
Dehydrated	0W	9.7 ... 10.2
Monohydrated	1W	11.6 ... 12.9
Bihydrated	2W	14.9 ... 15.7
Trihydrated	3W	18 ... 19

185 However, even under controlled conditions different hydration states coexist and are  
186 common in smectites <sup>[53]</sup>. Hence, transition states, such as 0W/1W, 1W/2W and 2W/3W, and  
187 even the coexistence of three hydration states are possible. Furthermore part of the  
188 interlayer water is connected to interlayer cations and part is connected to the interlayer clay  
189 mineral surface<sup>[54]</sup>. The water binding mechanisms on clay minerals are complex.  
190 Nevertheless, the increase of the  $d001$  is useful to display the expansion of the interlayer due  
191 to water uptake between the clay sheets. In the following, the  $d001$  will be used to  
192 demonstrate the swelling of the interlayer caused by the gradual water uptake for reasons of  
193 simplifying this complex topic.

194

#### 195 *Determination of water content*

196 The moisture of the samples was determined by static and dynamic heating. Static heating  
197 was performed in duplicate by heating about 500 mg of the sample in glasses for 24 h at  
198 105 °C and determining the weight loss caused by heating <sup>[55]</sup>. Hereby, the geotechnical

199 water content ( $W_{gt}$ ) was determined (eq. (1)).  $W_{gt}$  refers to the dry weight of the sample as  
 200 the reference base and is calculated as the ratio of mass of water ( $m_w$ ) to dry weight ( $m_d$ ) of  
 201 the sample.

202

$$W_{gt}[\%] = \frac{m_w}{m_d} \cdot 100 \quad (1)$$

203

204 Dynamic heating was realized with Simultaneous Thermal Analysis (STA), where the thermal  
 205 properties were determined in a defined temperature-time-program with the device  
 206 STA 449 C Jupiter (Netzsch, Germany) and the experimental parameters listed in Table 3.

207

208 Table 3: Experimental parameters for STA measurements.

Parameter	Agent	Amount	Unit
crucible	Pt/Rh, with lid		
reference	empty Pt/Rh crucible with lid		
purging gas	synthetic air	50	ml/min
protective gas	nitrogen, quality 6.0 (99.999%)	20	ml/min
start temperature	-	35	°C
final temperature	-	1100	°C
heating rate	-	10	K/min
net weight	sample	100	mg

209

210 The mass loss of the dehydration was determined from the Thermogravimetry curve (TG)  
 211 and was used to determine  $W_{gt}$ . With help of Differential Scanning Calorimetry (DSC) and the  
 212 water release observed from Evolved Gas Analysis (EGA) with a coupled mass  
 213 spectrometer, a material-specific temperature range of dehydration was determined, which  
 214 was in the range of 35 – 250 °C for ben\_Cal and ben\_P, and 35 – 200 °C for ben\_Vol.

215

216

### 217 *Dielectric measurements*

218 The dielectric properties were studied with the ALPHA-AN-Analyzer (Novocontrol) and the  
 219 software WinDETA 5.0 within a frequency range of  $10^{-4}$  to  $10^6$  Hz and a fixed voltage of

220 1.5 V. After a small sample amount of about 500 mg was mounted between the two  
 221 electrode plates made of stainless steel, the plates were adjusted with a micrometer screw  
 222 and the thickness noted. The upper electrode incorporated a guard ring in order to eliminate  
 223 edge and surface current effects, such as stray capacitance at the electrode edges as well  
 224 as leakage current flows on the sample surface. The upper signal electrode had a diameter  
 225 of 10 mm and the lower electrode a diameter of 30 mm. Thus, the effective sample volume  
 226 was determined by the upper signal electrode incorporating the guard ring (Figure 1). After  
 227 mounting the sample, the sample cell was transferred into a stainless steel container, which  
 228 was prepared with the same salt solution that was used for moisture pre-equilibration of the  
 229 material. Thereafter, the container was sealed with Parafilm<sup>®</sup> (Merck) for the duration of the  
 230 measurement to avoid humidity influences from the environment. The sample cell used is  
 231 described in further detail elsewhere <sup>[56]</sup>.

232

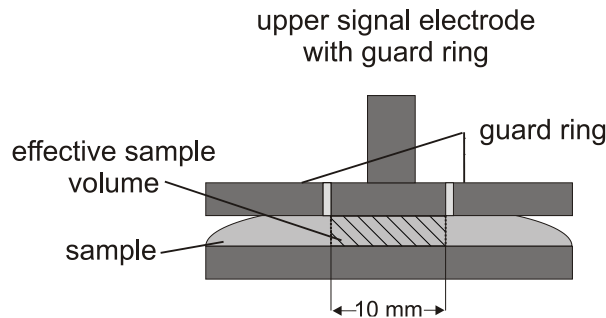


Figure 1: Schematic set up of the electrode plates with mounted sample.

233 For each of the bentonites, complex permittivity, complex resistivity and complex electrical  
 234 conductivity were determined. The relationship between these parameters is shown in eq.  
 235 (2), (7) and (5) below. The complex dielectric permittivity ( $\epsilon^*$ ) depends on angular frequency  
 236 ( $\omega = 2\pi f$ ) and can directly be calculated from the complex capacitance  $C^*$ .

$$\epsilon^*(\omega) = \frac{C^*(\omega)}{C_0}, \quad (2)$$

237 where  $C_0 = \epsilon_{air} A/d$  is the capacitance of the empty cell, A is the electrode area and d is sample  
 238 thickness. The real part of the dielectric permittivity is denoted by  $\epsilon'$  and the imaginary part

239 by  $\varepsilon''$ . The imaginary part of permittivity characterizes the dielectric loss of a medium and  
240 comprises contributions from relaxation processes ( $\varepsilon''_{relax}$ ) and from the dc conductivity ( $\sigma_{dc}$ ):

$$\varepsilon^* = \varepsilon' - i \left( \varepsilon''_{relax} + \frac{\sigma_{dc}}{\omega \varepsilon_0} \right) \quad (3)$$

241 where  $\varepsilon_0$  describes the permittivity of vacuum being  $\sim 8.85 \cdot 10^{-12}$  F/m.

242 The dissipation factor ( $\tan \delta$ ), which can be used to identify relaxation processes, is given by  
243 eq. (4):

$$\tan \delta = \frac{\varepsilon''}{\varepsilon'} \quad (4)$$

244

245 The complex conductivity  $\sigma^*$  is related to the complex permittivity as

$$\sigma^* = i\omega \varepsilon_0 \varepsilon^*, \quad (5)$$

246 and it comprises a real ( $\sigma'$ ) and an imaginary part ( $\sigma''$ ):

$$\sigma^*(\omega) = \sigma'(\omega) + i\sigma''(\omega) \quad (6)$$

247

248 Further, the relation between the complex resistivity and the complex conductivity is  
249 expressed as

$$\rho^*(\omega) = \frac{1}{\sigma^*(\omega)}. \quad (7)$$

250

### 251 *Equivalent circuit modelling*

252 Equivalent circuit modelling was implemented using the software ZView (version 3.2c,  
253 Scribner Associates) in order to be able to describe the response of the bulk system by  
254 idealized model circuits. The quality of the fit describing the match of the sample impedance  
255 with that of the modelled impedance of the equivalent circuit is described by the weighted  
256 sum of squares ( $\chi^2$ ).

257

258

259 RESULTS AND DISCUSSION

260 *Moisture / adsorbed water states*

261 The hydration of the bentonites with increasing r.h. could clearly be observed from  
 262 gravimetric water contents determined by static and dynamic heating (Table 4). The  
 263 maximum amount of adsorbed water at 93% r.h. increased with increasing smectite content  
 264 and increasing CEC (ben\_P > ben\_Vol > ben\_Cal). However, this dependency was different  
 265 for the samples stored at 11% and 43% r.h.. For these lower relative humidities the amount  
 266 of adsorbed water sorted in the order ben\_P > ben\_Cal > ben\_Vol. This order equals the  
 267 order of the  $A_{s, out}$ . Hence, the water adsorption at low and intermediate r.h. was influenced  
 268 by multiple factors, such as porosity of the sample, CEC and type of interlayer cation with the  
 269 specific hydration energy, whereas the water adsorption at high r.h. was mainly influenced by  
 270 the CEC and  $A_{s, total}$ , i.e. the water adsorption at the inner surface for instance by hydration of  
 271 interlayer cations.

272

Table 4: Gravimetric water content determined by static and dynamic heating as well as hydration state of the smectites of the three bentonites stored at different relative humidities.

Sample	Parameter	Upper temperature		Humidity [%]		
				11	43	93
ben_Cal	static heating	105 °C	$W_{gt}$	7.3	9.4	18.5
	dynamic heating	250 °C	[%]	10.3	12.7	18.9
	hydration state	-	-	1W/2W	1W/2W	2W/3W
ben_Vol	static heating	105 °C	$W_{gt}$	4.4	7.3	19.6
	dynamic heating	200 °C	[%]	6.3	8.5	19.5
	hydration state	-	-	0W/1W	1W	2W
ben_P	static heating	105 °C	$W_{gt}$	8.0	13.3	28.9
	dynamic heating	250 °C	[%]	15.6	17.0	30.4
	hydration state	-	-	1W/2W	1W/2W	2W

273

274 The water content determined by dynamic heating was higher than the water content  
 275 determined by static heating at 105 °C. This implies an incomplete removal of adsorbed  
 276 water by static heating for smectite-rich clays with the commonly applied reference method.  
 277 The underestimation of static  $W_{gt}$  in comparison to dynamic  $W_{gt}$  was most severe at low r.h.  
 278 and reached up to 7.6% of absolute difference for ben\_P. This is most likely a result of the

279 higher binding energy of the water molecules on the smectites at low relative humidity.  
 280 Complete removal of the strongly bound water required higher temperatures and therefore  
 281 resulted in a more severe underestimation of  $W_{gt}$  at 11% r.h. by static heating.  
 282 Due to their different interlayer occupation, the smectites of the bentonites showed different  
 283 hydration states with varying relative humidity (Figure 2). At 93% r.h. the smectites of all  
 284 three bentonites showed an expansion equalling a predominantly 2W hydration state and  
 285 ben\_Cal reached a transition state 2W/3W according to Table 2 and [53]. At 43% r.h. ben\_Cal  
 286 and ben\_P primarily showed a 1W/2W transition state, whereas the smectite of ben\_Vol  
 287 exhibited an expansion corresponding to 1W state. This behaviour displays the different  
 288 hydration properties of the interlayer cations (Table 1). Similarly, at low r.h. (11%) ben\_Cal  
 289 and ben\_P still showed a 1W/2W transition state, since these two contain the highest amount  
 290 of divalent cations within the interlayer and the hydration enthalpy of the cations decreases in  
 291 the order  $\text{Ca}^{2+}$  (-1920 kJ/mol) >  $\text{Mg}^{2+}$  (-1650 kJ/mol) >  $\text{Na}^+$  (-405 kJ/mol) [57]. Ben\_Vol, with a  
 292 dominance of monovalent cations showed a 0W/1W transition state at 11% r.h..

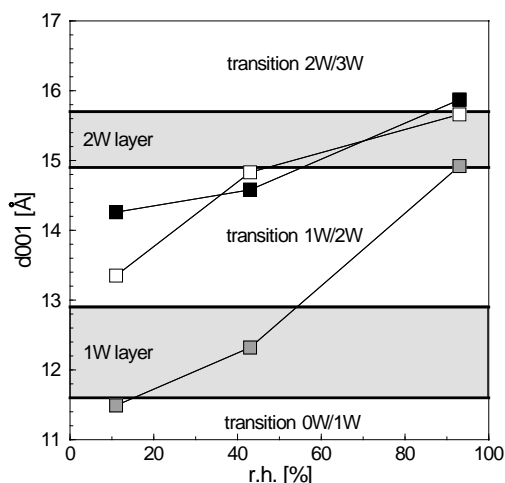


Figure 2: Basal spacing of the smectites in the bentonites in dependence of relative humidity. ben\_Cal (■), ben\_Vol (◐), ben\_P (□)

293

### 294 Dielectric spectra

295 The difference in hydration state and  $W_{gt}$  as a function of r.h. could furthermore be observed  
 296 in a changing dielectric behaviour. The real part of complex permittivity increases as a  
 297 function of water content over the whole frequency range (Figure 3).

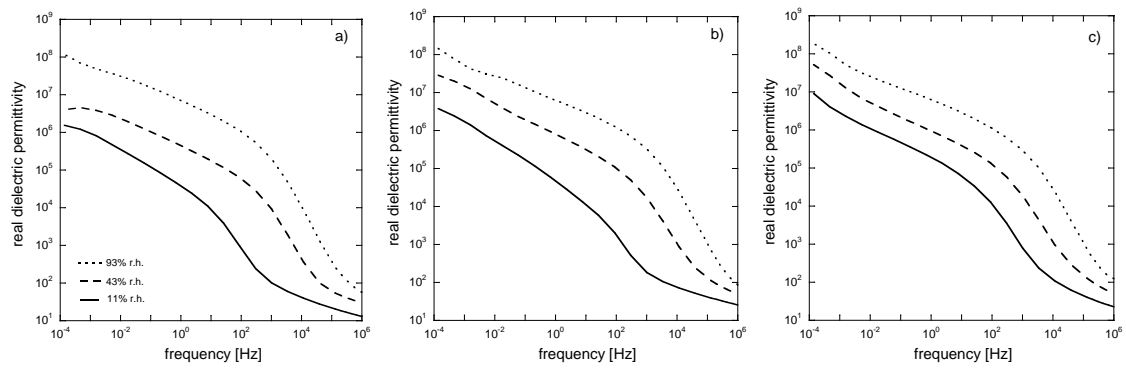


Figure 3: Real permittivity as a function of frequency for different r.h. a): ben\_Cal, b): ben\_Vol and c): ben\_P; the legend on the left applies for all graphs.

298

299 The regions of steepest slope in the real permittivity plots are assumed to correspond to a  
 300 relaxation in the imaginary part of complex permittivity, although due to superimposing  
 301 conductivity effects, no clear relaxation peaks could be identified in the imaginary complex  
 302 permittivity. Thus, the three bentonite samples showed a relaxation shift to higher  
 303 frequencies with increasing moisture, which was more clearly seen in the dissipation factor  
 304 as discussed next.

305 For each of the bentonites, the dissipation factor was higher the higher the water content and  
 306 the peaks shifted to higher frequencies (Figure 4). The peaks in the range of  $10^2$  to  $10^6$  Hz  
 307 are assumed to be due to interfacial processes such as polarization in the electrochemical  
 308 double-layer at the clay/water interfaces. Hence, interfacial processes were promoted by an  
 309 increased water content.

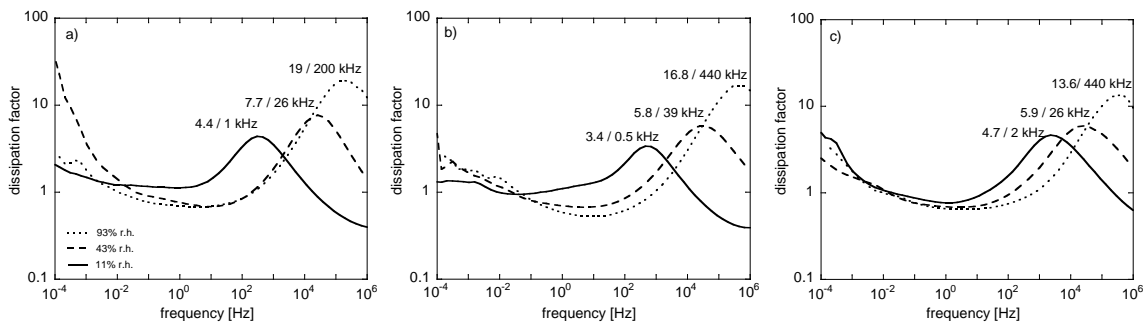


Figure 4: Dissipation factor as a function of frequency in dependence of water content; a): ben\_Cal, b): ben\_Vol and c): ben\_P; the legend on the left applies for all graphs.

310



311 The increase of the dissipation factor to values >1 at the lower end of the frequency range,  
 312 showing where the dielectric losses outweigh the dispersion, is likely due to electrode  
 313 polarization.

314 Similarly to the plots of real permittivity, the moisture difference could clearly be observed in  
 315 the real conductivity plots (Figure 5). A conductivity plateau at upper frequencies of the  
 316 measurement range was visible for each sample at all of the studied relative humidities, from  
 317 which the ionic conductivity of the clays could be obtained (Table 5). Compared to ionic  
 318 conductivity measured in the clear supernatant of centrifuged bentonite suspensions, the  
 319 ionic conductivity of the samples stored at 93% r.h. is of about the same order of magnitude  
 320 (data not shown).

321

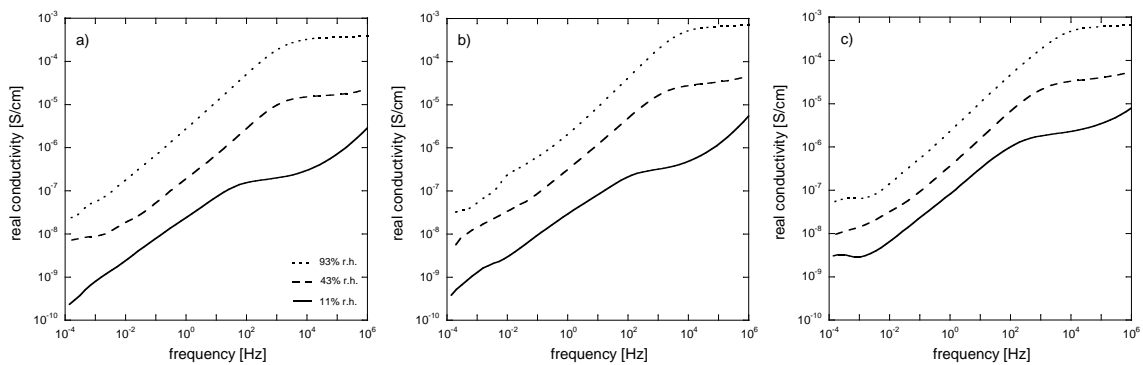


Figure 5: Real conductivity as a function of frequency for different r.h.; a): ben\_Cal, b): ben\_Vol and c): ben\_P; the legend on the left applies for all graphs.

322

Table 5: Ionic conductivity of the bentonites obtained from the conductivity plateau in dependence of equilibrated moisture conditions.

sample	r.h. [%]	$\sigma'$ [S/cm]
ben_Cal	11	1.5E-07
	43	1.3E-05
	93	3.0E-04
ben_Vol	11	1.8E-07
	43	2.4E-05
	93	5.7E-04
ben_P	11	1.7E-06
	43	2.9E-05
	93	5.2E-04

323

324 In all cases, the ionic conductivity of the bentonites increased with increasing water content,  
325 which is in agreement to the findings of Logsdon and Laird <sup>[58]</sup> and Bidadi *et al.* <sup>[59]</sup> who  
326 conducted dielectric studies of smectites at low water contents in the MHz range and on clay-  
327 films in the low frequency range, respectively.

328 Towards lower frequencies the conductivity plateau passed over to a linear section for each  
329 bentonite. The slope of the linear section was about 0.5, indicating diffusion at least as a  
330 participating underlying process in this frequency range.

331  
332 Electrode polarization (EP) occurred at the lower end of the frequency range and was  
333 already indicated by the dissipation factor. The occurrence of EP was proven experimentally  
334 by measuring on bentonites with different sample thicknesses. The permittivity is a bulk  
335 quantity and should be independent of thickness, while the capacitance  $C \sim 1/d$ , according to  
336 eq. (2). However, a spurious thickness dependence of the permittivity was observed at low  
337 frequencies, and this indicates the dominance of electrode polarisation at frequencies below  
338 1 Hz. Thus, the behavior at the lower end of the frequency range can be assumed to be due  
339 to both bulk response and electrode polarization, whereas the higher end of the frequency  
340 range can be assumed to be mostly due to bulk response (Figure 6).

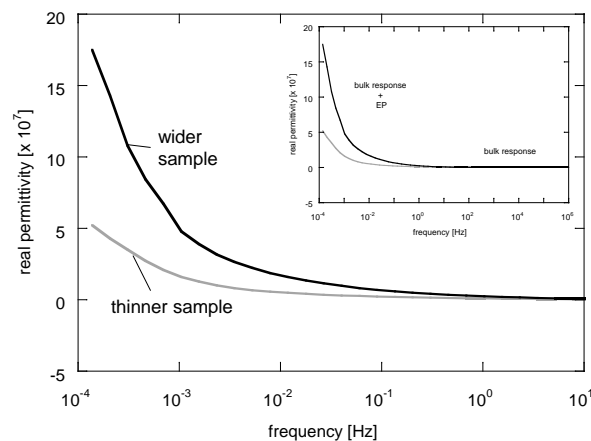


Figure 6: Real permittivity of ben\_P at 43% r.h. and different sample thicknesses to identify electrode polarization; black curves: higher sample thickness (1.53 mm); grey curves: smaller sample thickness (0.97 mm).

341

342

343 *Equivalent circuit modelling*

344 Equivalent circuit modelling was frequently used in the past to identify the mechanisms  
345 taking place in clay suspensions and clay-water electrolyte systems, including smectitic  
346 clays, at low frequencies [40, 48]. Although some models were adjusted to describe double-  
347 layer polarization and Maxwell-Wagner polarization, and accounted for the accumulation of  
348 charges on the electrode due to electrode polarization, they showed to be completely  
349 inappropriate for the three presently studied bentonites at low moisture level. Hence, different  
350 dielectric mechanisms take place in clays dominated by free water rather than bound water  
351 and new equivalent circuit models are necessary to describe the effects taking place.

352 The following equivalent circuit model was chosen to describe the measured impedance and  
353 it is based on considerations of the physics of relaxation processes in a humid porous  
354 material. (Figure 7). The circuit contains constant phase elements (CPE) and resistances to  
355 account for bulk conductivity and electrochemical double-layer polarization.

356

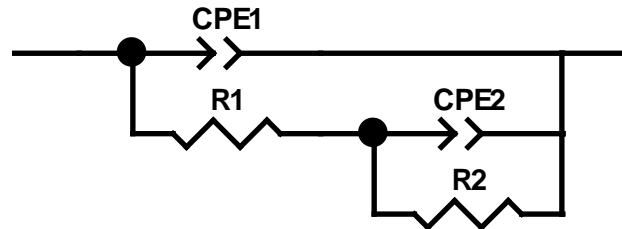


Figure 7: Equivalent circuit model representing the clay-water-mixture, consisting of resistances R1 and R2 and constant phase elements CPE1 and CPE2.

357

358 The CPEs were used instead of capacitors to account for heterogeneities in the relaxation at  
359 the clay-water-interface. The CPE impedance ( $Z_{CPE}$ ) is characterized by the parameters  $T$   
360 and  $P$  (eq. (8));

361

$$Z_{CPE} = \frac{1}{T(i\omega)^P}, \quad (8)$$

362

363 where  $T = (\tau^{*P})/B$ ,  $B$  is the amplitude of  $Z_{CPE}$ ,  $\tau^*$  is a characteristic time constant and  $P$  an  
364 exponent varying between zero and one that describes the constant phase angle in the  
365 complex plane plot. A  $45^\circ$  line, signifying a diffusion process, is produced in the complex  
366 plane with  $P = 0.5$ . For  $P = 1$ , the CPE impedance equals that of a capacitor and for  $P = 0$   
367 that of a resistance.

368 The element CPE1 is attributed to the capacitive part of a high frequency relaxation, which  
369 can be due to bound water or to the Maxwell-Wagner effect <sup>[60]</sup>. Putting an additional  
370 resistance in series with CPE1 to model a high frequency relaxation did not improve the fit,  
371 and resulted in very high statistical errors in the value of the added series resistance. The  
372 resistance R1 is associated to the movement of ions dissolved in the adsorbed water in the  
373 pores of the bentonites. In addition, the series connection of R1 and CPE2 can be interpreted  
374 as a model of the electrochemical double-layer relaxation in the pore fluid close to the  
375 interfaces. Charges at the surface of the solid attract counter-ions of opposite charge and  
376 these give rise to the Stern layer. At a larger distance from the interface a diffuse layer  
377 develops at sufficiently high water content. The counter-ion concentration decreases with  
378 distance from the interface and the concentration of the oppositely charged ions increase.  
379 The double-layer relaxation of a sphere in a fluid has been studied frequently and a review of  
380 a number of approaches can be found in Nettelblad and Niklasson <sup>[61]</sup>. The most complete  
381 theory appears to be that of DeLacey and White <sup>[62]</sup>. It was found by Nettelblad and Niklasson  
382 <sup>[61]</sup> that a simple phenomenological expression of the Havriliak-Negami (HN <sup>[63]</sup>) type could  
383 give satisfactory fits to numerical results of the dielectric permittivity for a number of theories  
384 for the double-layer relaxation at a dielectric sphere. Hence we use the expression

385

$$C_{HN}(\omega) = \frac{A}{[1 + (i\omega\tau)^{1/2}]^2} \quad (9)$$

386

387 to represent the frequency-dependent capacitance of the double-layer relaxation. Here  $A$  is  
388 the relaxation strength and  $\tau$  is the relaxation time. However, a porous material is not

389 composed of solid spheres, and in a porous structure the solid particles are connected to one  
 390 another.

391  
 392 The dielectric response of a water-filled pore can be described by transmission line models  
 393 [64]. The simplest model for a uniform pore assumes a distributed resistance along the pore  
 394 and a distributed capacitance at the pore interface. A number of generalizations taking into  
 395 account anomalous diffusion processes as well as trapping of charges were developed by  
 396 Bisquert *et al.* [65, 66]. We then combined the transmission line approach with eq. (9) and took  
 397 into account both the electrochemical double-layer relaxation and the conductivity due to ions  
 398 dissolved in the pore fluid. The ion conduction process in the pore is represented by the  
 399 resistances (R) in the upper part of Figure 8, while the double-layer relaxation is represented  
 400 by the HN impedance connecting the resistance line with the pore interface.

401

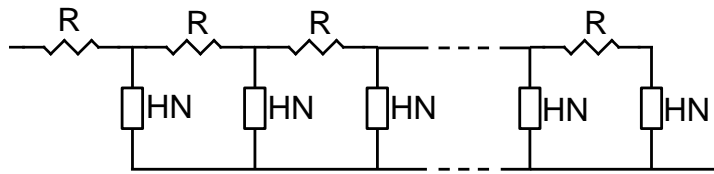


Figure 8: Transmission line for the double-layer relaxation in a single pore. The resistances describe ion conduction in the pore space, while HN is a Havriliak-Negami impedance representing the double-layer relaxation.

402

403 Converting the HN equation (eq. (9)) to impedance we find,

$$Z_{HN}(\omega) = \frac{[1+(i\omega\tau)^{1/2}]^2}{i\omega A}. \quad (10)$$

404

405 A transmission line has also a termination at the right hand end, which has not been  
 406 specified in Figure 8. Considering the bulk part of the dielectric response and neglecting  
 407 termination effects, gives the following impedance for the transmission line in Figure 8:

$$Z(\omega) = (R \cdot Z_{HN}(\omega))^{1/2} = \frac{R^{1/2}(1 + (i\omega\tau)^{1/2})}{(i\omega A)^{1/2}}. \quad (11)$$

408

409 It is easily found that a series R-CPE circuit has an impedance

$$Z(\omega) = \frac{B+R(i\omega\tau^*)^P}{(i\omega\tau^*)^P}, \quad (12)$$

410

411 which is found to be of the same form as equation (11) when the power-law exponent  $P = 0.5$

412 The power law exponent of the double-layer relaxation may depart from the value 0.5 for a

413 variety of reasons, such as branched transmission line networks <sup>[67]</sup> or rough and fractal pore

414 surfaces <sup>[68]</sup>. We conclude that the occurrence of a series R-CPE combination in an

415 equivalent circuit analysis of experimental data can be interpreted in terms of a double-layer

416 relaxation. The interpretation of the series resistance R2 is more uncertain but we presume

417 that it is related to the termination of the transmission line. This resistance is important only

418 at the low frequency end of the spectra at  $< 1$  Hz.

419

420 The fit parameters of the circuit described in Figure 7 are listed in Table 6. The model gives a

421 satisfactory description of the measured values (Figure 9) with low error estimates and sum

422 of  $X^2$  ranging between 0.23 and 0.75 for all three bentonites.

423

Table 6: Fit parameters of the equivalent circuit model accounting for bulk effects and interfacial effects in clays at low moisture level. The fit parameters are given as resistivities/permittivities and are thus independent of sample geometry.

	r.h. [%]	sum $X^2$	CPE1-T [F/m*s <sup>(1-P)</sup> ]	CPE1-P	R1 [ohm*m]	CPE2-T [F/m*s <sup>(1-P)</sup> ]	CPE2-P	R2 [ohm*m]
ben_Cal	11	0.53	6.89E-10	0.74	487740	1.30E-07	0.55	8.72E+08
	43	0.51	4.82E-10	0.82	6801	7.51E-07	0.63	1.40E+07
	93	0.50	4.10E-08	0.56	273	1.29E-05	0.65	5.19E+06
ben_Vol	11	0.23	6.89E-10	0.78	269260	1.65E-07	0.50	1.49E+09
	43	0.75	2.59E-10	0.89	4066	1.36E-06	0.58	5.58E+07
	93	0.65	7.14E-07	0.43	133	1.22E-05	0.66	3.95E+06
ben_P	11	0.24	8.30E-09	0.64	34678	5.93E-07	0.61	3.27E+07
	43	0.29	1.12E-08	0.66	2941	1.93E-06	0.61	1.72E+07
	93	0.33	8.33E-08	0.57	180	9.90E-06	0.65	2.80E+06

424

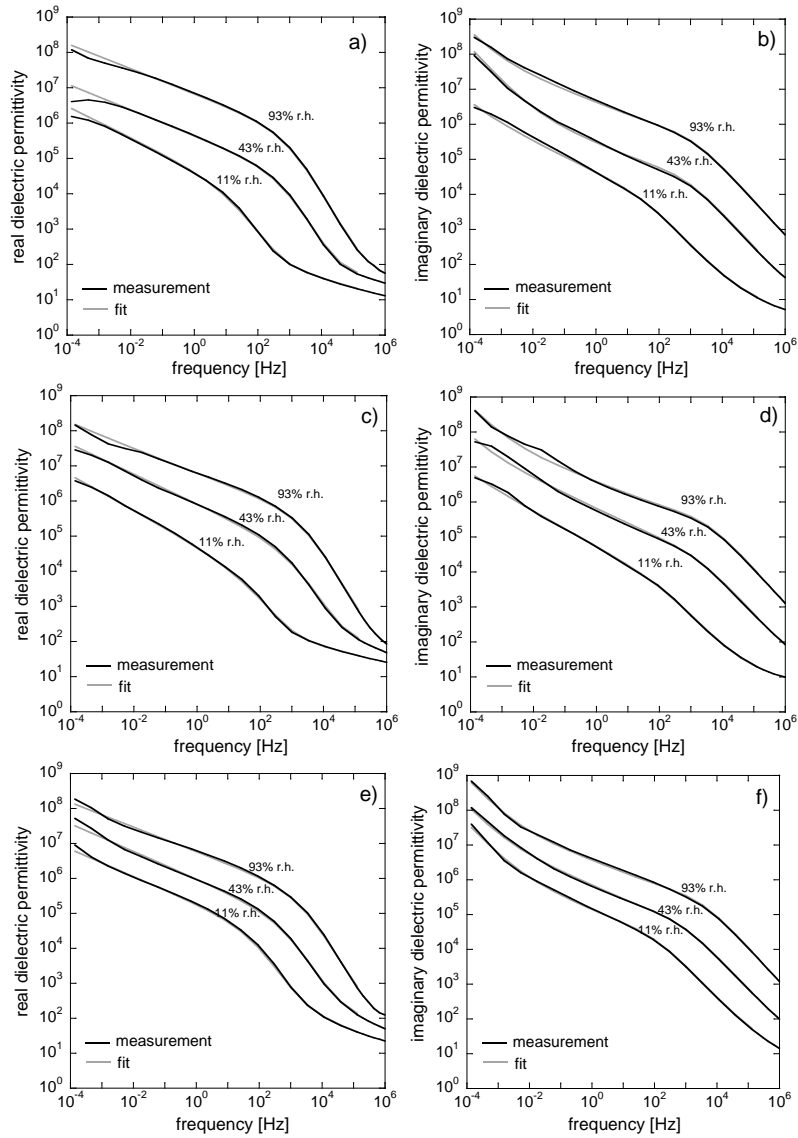


Figure 9: Fit results of the equivalent circuit described in Figure 7 at different r.h. conditions for ben\_Cal (a, b), ben\_Vol (c, d) and ben\_P (e, f).

425

426 A relationship between  $W_{gt}$  as well as r.h. for each discrete circuit element was observed  
 427 (Figure 10). The resistivity of both elements R1 and R2 decreased with increasing moisture,  
 428 which is due to a better conductivity of the samples at higher water content (Table 5). At 11%  
 429 and 43% r.h. the resistivity of R1 decreased in the order ben\_Cal < ben\_Vol < ben\_P  
 430 corresponding to the order of an increasing smectite content and increasing CEC. Hence, at  
 431 low water contents between 6 and 17% STA  $W_{gt}$ , a high amount of swellable 2:1 layer  
 432 silicates and a high ability to reversibly exchange cations resulted in a reduced bulk  
 433 resistivity. This implies that the Cu-Trien exchangeable cations, preferably those bound to  
 434 the smectitic surface, are responsible for the conductivity at least at water contents up to

435 17%  $W_{gt}$ . At those low water contents the interlayer cations are possibly as strongly bound as  
436 the cations associated with the outer clay mineral surface forming the Stern layer.  
437 At water contents of 19% STA  $W_{gt}$  and above other parameters than smectite content and  
438 CEC seemed to become more important. At 93% r.h. the resistivity of R1 increases in the  
439 order  $benVol < benP < benCal$  which corresponds to the order of the basal spacing  $d001$   
440 and to the amount of divalent cations within the interlayer (Table 1 and Table 2).  
441 This coincides with findings in the literature <sup>[10, 69]</sup> with slightly different thresholds. Herein it  
442 was concluded that the reorientation of water molecules bound to clays with a water content  
443 of up to 4%  $W_{gt}$  takes significantly longer than the reorientation of water molecules bound to  
444 clays with water contents of about 13%  $W_{gt}$ . Hereby, a longer reorientation time of bound  
445 water molecules was assumed to be associated with a stronger binding of the water  
446 molecules to the interlayer cation of the clay mineral. Further, the reorientation of the water  
447 molecules bound to the interlayer cation showed to be different from that of remaining water  
448 <sup>[10, 70]</sup>. However, further analysis, for instance with homoionically exchanged samples, are  
449 necessary to clearly attribute the behavior to individual processes.  
450



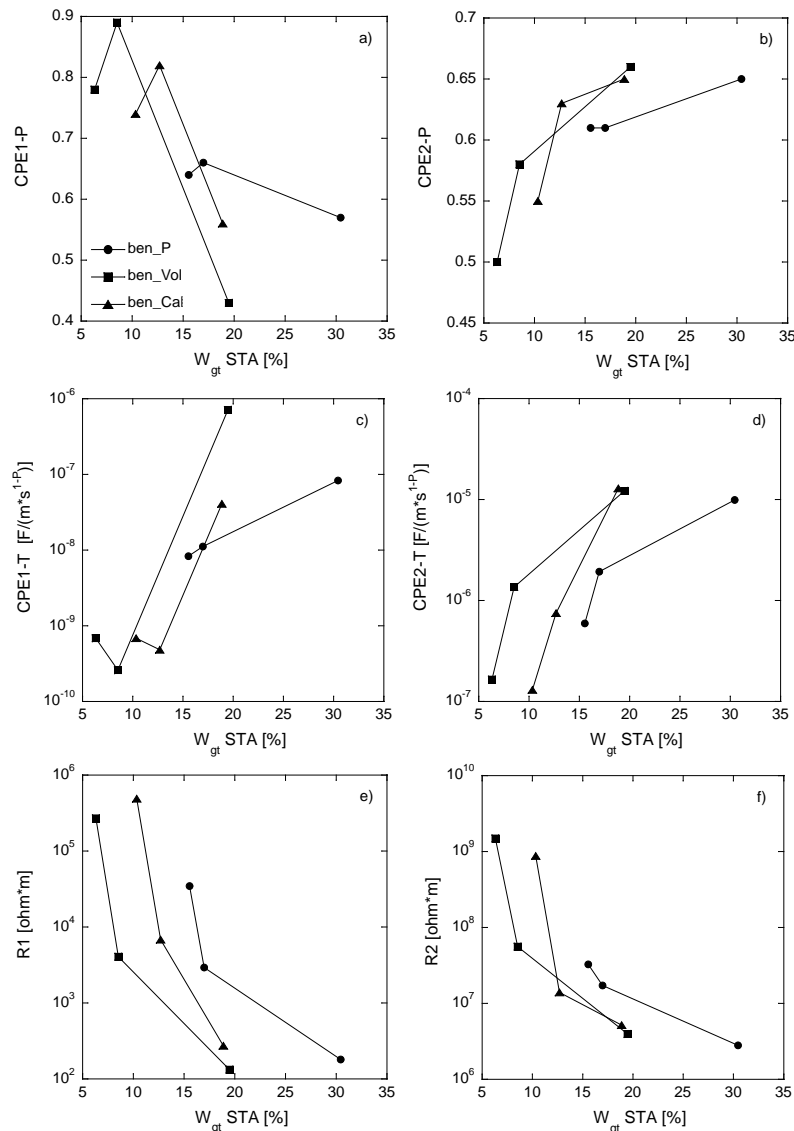


Figure 10: Relationship of discrete circuit elements and STA  $W_{gt}$  for the equivalent circuit model of Figure 7; the legend in plot a) applies for all graphs.

451

452 The corresponding capacitive element CPE1 of the bulk response showed a similar behavior  
 453 of ben\_Vol and ben\_Cal, whereas ben\_P seemed to be characterized by partly different  
 454 processes (Figure 10).

455 The exponent  $P$  of CPE1, describing the constant phase angle in the complex plane plot,  
 456 was always highest at intermediate relative humidities and increased in the order  
 457 93% r.h. < 11% r.h. < 43% r.h. (Table 6, Figure 10). At 93% r.h. the CPE1  $P$  values ranged  
 458 between 0.43 and 0.56 and at 43% r.h. between 0.66 and 0.89. For ben\_P the constant  
 459 phase angle showed little variation (0.57 ... 0.66) between the studied moisture states. A  $P$   
 460 value of about 0.5 suggests diffusion as a contributing mechanism. At very low relative

461 humidities the *CPE2 P* values that characterize the interfacial double-layer relaxation  
462 process, suggest a possible occurrence of diffusion for the two bentonites ben\_Cal and  
463 ben\_Vol (0.55 and 0.5, respectively). With increasing relative humidity the exponent departs  
464 more and more from the value 0.5 .

465 The occurrence of electrode polarization was proven for ben\_P (Figure 6) and is also likely to  
466 be present in ben\_Vol and ben\_Cal at low frequencies. In order to assess the thickness  
467 dependency of the discrete circuit elements, the model of Figure 7 was applied to ben\_P at  
468 43% r.h. for different sample thicknesses (1.53 mm vs. 0.97 mm). Equivalent circuit modeling  
469 on samples of different thickness was performed on impedance level, hence without  
470 correction for area and thickness of the sample. Hereby, the circuit element R1 was larger for  
471 the thicker sample, which supports the interpretation that it describes the bulk resistivity. In  
472 addition, CPE1-T was characterized by almost the same thickness dependency as R1,  
473 strongly indicating that these elements are coupled and represent a single mechanism. The  
474 other resistance R2 influenced the spectra below 0.1 Hz for ben\_P at 43% r.h. It showed  
475 lower values for thicker samples, i.e. the opposite behavior to R1. This behavior is clearly not  
476 a bulk one and is presumed to be due to a termination of the transmission line or to charge  
477 transfer mechanisms at the electrodes.

478 In a second model (Figure 11) a separate capacitor was added in series with the other  
479 discrete circuit elements in order to separate a possible capacitive part of the electrode  
480 polarization effect from other processes. It led to better fits at intermediate and high water  
481 contents for ben\_P, as seen in Table 7.

482

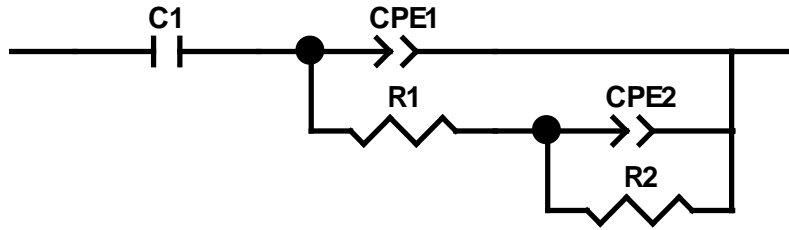


Figure 11: Equivalent circuit model for a bentonite of high purity (ben\_P) that separately accounts for capacitive effects of electrode polarization.

483

484

Table 7: Fit parameters of ben\_P for the model given in Figure 11.

r.h.	sum $X^2$	C1	CPE1-T	CPE1-P	R1	CPE2-T	CPE2-P	R2
[%]		[F]	[F/s <sup>(1-P)</sup> ]		[ohm]	[F/s <sup>(1-P)</sup> ]		[ohm]
43	0.0797	0.0006	3.31E-07	0.44	3480	1.49E-06	0.65	1.38E+07
93	0.2894	0.0073	6.31E-08	0.60	181	9.72E-06	0.65	2.50E+06

485

486 The second model (Figure 11) depended on the smectite content of the bentonite and the  
 487 water content. It worked reliably at intermediate and high relative humidities for high smectite  
 488 contents. At 11% r.h. and for the two bentonites with a higher amount of impurities, the  
 489 model was characterized by high errors and high sum of  $X^2$  values.

490 The build-up of charges on the electrode was higher at higher water content, as could be  
 491 described by the capacitor C1 (Table 7). The accumulated charges gave rise to capacitance  
 492 values in the range of  $0.6 \cdot 10^{-3}$  and  $7.3 \cdot 10^{-3}$  F for ben\_P.

493 Nevertheless, it can be expected that electrode polarization exists for the bentonites with  
 494 lower smectite content and possibly for ben\_P at 11% r.h.. However, then EP could not  
 495 clearly be separated from the other processes in the equivalent circuit.

496

497

498

499 SUMMARY AND CONCLUSIONS

500 The low frequency dielectric properties of three bentonites with varying smectite content  
501 were analyzed after storage at defined relative humidity. The clays exhibited low water  
502 contents ranging between 6 and 30% STA  $W_{gt}$ .

503 Dielectric properties were clearly influenced by hydration state and water content. Real  
504 dielectric permittivity and real conductivity increased with increasing amount of adsorbed  
505 water. The imaginary dielectric permittivity showed no clear relaxation peak due to  
506 superimposing conductivity effects. However, the dissipation factor  $\tan \delta$  revealed clear  
507 peaks for each of the bentonites in the range of  $10^2$  to  $10^6$  Hz that were attributed to  
508 interfacial polarization such as polarization in the electrochemical double layer at the clay-  
509 water interface. The intensity of the interfacial processes were promoted by increased water  
510 content and shifted to higher frequencies.

511 The occurrence of electrode polarization was shown to be a participating process at the  
512 lower end of the frequency range ( $\sim < 1$  Hz). However, separating electrode polarization from  
513 other processes in equivalent circuit modeling could only be achieved for the purest of the  
514 bentonites (ben\_P) with a smectite content of 98%. The charges that accumulated on the  
515 electrode gave rise to capacitance values of about  $0.6 \cdot 10^{-3}$  and  $7.3 \cdot 10^{-3}$  F.

516 An equivalent circuit model consisting of a combination of CPEs and resistances was found  
517 to successfully describe the low frequency dielectric behavior of bentonites at adsorbed  
518 water states. A series R-CPE connection was used to model the electrochemical double-  
519 layer relaxation in the pore fluid close to the interfaces. Equivalent circuit modelling further  
520 revealed that bulk resistivity was mainly influenced by the smectite content and CEC at water  
521 contents up to 17% STA  $W_{gt}$ . At water contents of  $\geq 19\%$  STA  $W_{gt}$  structural parameters of  
522 the interlayer, such as interlayer occupation and clay mineral swelling due to water uptake  
523 into the interlayer predominantly influenced the bulk resistivity.

524 In order to validate our model on a larger set of materials and ascertain the importance of  
525 different structural and mineralogical parameters on the dielectric response, further dielectric

526 analysis on clays at low moisture levels, especially homoionic exchanged bentonites, are  
527 necessary and currently performed.

528

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536

537

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